

Figure 1. (a) Room temperature optical absorption spectra of metal-12-crown-4 solutions; (b) room temperature pulse radiolysis spectrum (arbitrary units) of neat 12-crown-4

here are in fair agreement with earlier reports of Na^- ($\lambda_{\text{max.}} = 660 \pm 20 \text{ nm}$) and K^- ($\lambda_{\text{max.}} = 890 \pm 20 \text{ nm}$) in 12-crown-4.⁹ Note also the presence of Na^- in K, Rb, and Cs solutions, almost certainly originating from sodium extracted from the Pyrex preparation cell.

No i.r. absorptions which could be attributed to solvated electrons were detected in the metal solutions, despite an earlier report of a band due to e_s^- in K-12-crown-4 at 1 800 nm.⁹ The optical spectrum of e_s^- in 12-crown-4, obtained by pulse radiolysis of the pure solvent, is shown in Figure 1(b). The end of pulse spectrum (0.2- μs pulses) shows a peak between wavelengths of 1 100 and 1 500 nm; assignment to e_s^- was based on the complete disappearance of this peak when the solution was saturated with the electron scavenger N_2O .^{*} The difference in absorption wavelength between the K-12-crown-4⁹ and pulse radiolysis results may be an effect of metal concentration; a similar phenomenon is reported in metal-ammonia solutions, where $\lambda_{\text{max.}}$ shifts to longer wavelength with increasing concentration.¹⁰ Alternatively, the band at 1 800 nm⁹ may be a feature of the crown ether molecule itself; we have observed similar weak absorptions at 1 720 nm in the spectrum of pure 12-crown-4.

The wavelength of the solvated electron absorption maximum in 12-crown-4 is similar to that in ammonia (1 504 nm).¹¹ While $\lambda_{\text{max.}}$ strictly indicates the $1s \rightarrow 2p$ energy separation between the ground and excited states of e_s^- , a red shift is generally attributed to a decrease in electron solvation energy. Thus the chemical potential of e_s^- in 12-crown-4 is about the same as that in NH_3 .

In summary, the optical spectra reveal high concentrations of M^- in alkali metal-12-crown-4 solutions. Conversely, they provide no evidence for high concentration of e_s^- in these metal solutions. This is in accord with earlier work showing that in less polar solvents such as polyethers the ratio of metal anions to solvated electrons is generally larger than in liquid ammonia.⁴

N.M.R. Spectra.—The steady-state optical spectra do not allow distinctions to be made between the various possible structures for M^- in solution.³ The most logical description of this species is a 'genuine' alkali metal anion with two electrons in

the outer s orbital. However, the optical spectra could also be assigned to species such as a triple ion, $e_s^- \text{M}_s^+ e_s^-$, formed by aggregation of two solvated electrons and the solvated cation, or by an ion pair between M_s^+ and the electron pair e_2^{2-} , or by an 'expanded' anion in which two electrons are in an expanded orbital around a solvated metal cation.¹² The important feature of all these structures except the alkali metal anion is the presence of a solvated cation.

The observation of the ^{23}Na n.m.r. spectrum of Na^- in Na-12-crown-4 solutions provided conclusive evidence that Na^- is a genuine anion with two electrons in the $3s$ orbital.¹³ The n.m.r. spectrum of Na^- in 12-crown-4 is shifted little ($\delta = -61.8 \pm 0.1$ p.p.m. at 300 K relative to a $\text{Na}^+ \text{-D}_2\text{O}$ reference sample) from the value calculated for gaseous Na^- ($\delta = -63.1$ p.p.m.).¹⁴ The absence of any large paramagnetic shift and the extreme narrowness of the line (width at half-height, $\Delta\nu_{1/2} = 10$ Hz) show that the solvent is effectively excluded from interaction with the $2p$ electrons. This indicates that Na^- is an alkali metal anion with the $3s$ orbital screening the $2p$ electrons from interaction with the solvent and providing a highly symmetric environment. All other models for the species of stoichiometry M^- involve cation solvation and would introduce a solvent-dependent paramagnetic shift.

The ^{87}Rb n.m.r. spectrum of Rb^- in Rb-12-crown-4 solutions has been reported previously.¹³ For this anion a moderate paramagnetic shift from the value for Rb^- in the gas phase is observed: $\delta = -191 \pm 2$ p.p.m. for Rb^- in 12-crown-4 at 300 K, compared with the calculated value of -213.6 p.p.m. for gaseous Rb^- .¹³ In addition, the n.m.r. linewidth of Rb^- in 12-crown-4 is broad ($\Delta\nu_{1/2} = 1 000$ Hz at 300 K) and the nuclear relaxation rate is much higher ($T_1^{-1} \geq 1.25 \times 10^3 \text{ s}^{-1}$ at 297 K) than for Na^- in the same solvent ($T_1^{-1} = 5.35 \text{ s}^{-1}$ for a Na-K-12-crown-4 solution at 300 K). However, the assignment of the ^{87}Rb n.m.r. resonance to the rubidide ion is unambiguous due to the extremely large diamagnetic shift from the $\text{Rb}^+ \text{-D}_2\text{O}$ reference. The other experimental parameters suggest that the nuclear shielding of Rb^- in the liquid state is significantly decreased and that a considerable degree of interaction with the solvent exists for Rb^- . Incomplete shielding of the outer p electrons of Rb^- and Cs^- in metal-cryptand-solvent systems has been noted previously.¹⁴ Evidence for a measurable degree of solvent participation in the ground-state wavefunction of K^- and Rb^- has also come from a study of the photodetachment of electrons from solvated alkali metal anions by e.s.r.¹⁵

Although the optical spectra indicate the presence of high concentrations of potassium and caesium anions in 12-crown-4

* Extensive measurements will be reported elsewhere (A. Ellaboudy, D. M. Holton, J. Langan, G. A. Salmon, and P. P. Edwards, in preparation).

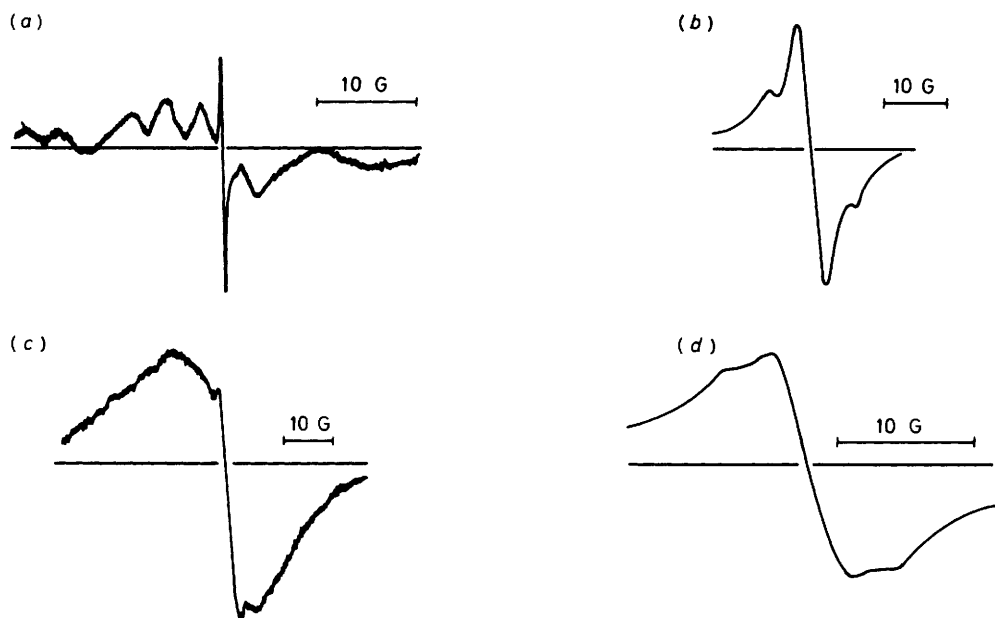
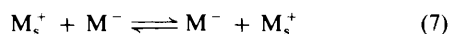
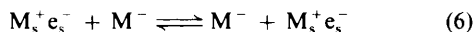


Figure 2. E.s.r. spectra of frozen M-12-crown-4 solutions: (a) M = Na, at 40 K, 1 mW microwave power; (b) M = K, at 120 K, 10 mW; (c) M = Rb at 125 K, 1 mW; (d) M = Cs at 25 K, 20 mW

solutions, we have been unable to detect these species by n.m.r. We stress that this is not a problem of sample instability. Exchange processes such as (6) and (7) which effectively limit



the lifetime of M^- in solution may be fast on the n.m.r. time-scale for K^- and Cs^- , causing the signals to be paramagnetically shifted and broadened beyond the limits of detection.

E.S.R. Spectra.—Metal-crown ether solutions are sufficiently stable in the liquid state for comparative e.s.r. measurements on liquid and solid phases to be made. Figure 2 shows the e.s.r. spectra obtained from frozen solutions of sodium, potassium, rubidium, and caesium in 12-crown-4. In each case, two paramagnetic species were detected which can be identified as the solvated electron, e_s^- , and the corresponding ion pair species, $M_s^+ e_s^-$. The solid phase spectra were not generally reproducible after thawing and refreezing. Fluid solutions exhibited only a single, time-averaged e.s.r. signal. It is interesting to compare the relative intensities of this signal for the various metals. Sodium and rubidium samples, in which Na^- and Rb^- have been detected by n.m.r., give e.s.r. signals approximately 100 and 10 times weaker respectively than the corresponding K and Cs spectra. Potassium samples provided the most comprehensive set of results and these are discussed in most detail.

(i) *Potassium.* Frozen K-12-crown-4 solutions (Figure 2) showed two superimposed e.s.r. signals at $g = 2.0020$ with peak-to-peak derivative linewidths $\Delta H_{pp} = 2.5 \pm 0.5$ and 12 ± 1 G at 120 K. The narrow and broad resonances are assigned respectively to the electron (e_s^-) and the ion pair ($K_s^+ e_s^-$). The broad linewidth of the ion-pair signal and its inhomogeneous 'power saturation behaviour' (discussed below) indicate unresolved hyperfine coupling to the metal nucleus. The ion-pair linewidth remained constant up to the melting point of the solution (ca. 290 K) while that of the electron decreased slightly to $\Delta H_{pp} = 2$ G near the melting temperature. The intensity of both signals diminished with heating.

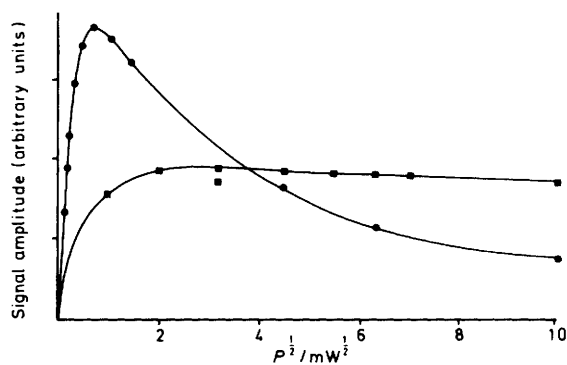


Figure 3. Power saturation of solvated electron (●) and ion-pair (■) resonances in frozen potassium-12-crown-4 solution at 120 K

The microwave power saturation behaviour of the two resonances is illustrated by Figure 3. The e_s^- resonance is a homogeneous signal; analysis of its saturation behaviour¹⁶ yields relaxation times $T_1 = (5 \pm 2) \times 10^{-5}$ s and $T_2 = (2.6 \pm 0.5) \times 10^{-8}$ s at 120 K; T_1 decreases with heating. The ion pair resonance has a Gaussian lineshape, and its power saturation (Figure 3) is typical of an inhomogeneous signal.¹⁶

On melting, the two resonances of frozen K-12-crown-4 solutions were replaced by a single, Lorentzian line at $g = 2.0020 \pm 0.0003$ with $\Delta H_{pp} = 0.06$ G. This resonance arises from rapid exchange between e_s^- and $M_s^+ e_s^-$; similar narrow Lorentzian signals are seen in liquid metal-ammonia¹⁷ and liquid metal-hexamethylphosphoramide (hmpa)¹⁸ solutions. The resonance is approximately ten times stronger than that from e_s^- in solid 12-crown-4 glasses just below the melting point. A power saturation study gave $T_1 = (5 \pm 1) \times 10^{-6}$ s, $T_2 = (1.1 \pm 0.2) \times 10^{-6}$ s at 297 K.

The characteristics of the resonances in metal-12-crown-4 solutions are summarised in the Table.

(ii) *Rubidium.* As indicated in Figure 2 and the Table, the e.s.r. characteristics of Rb-12-crown-4 solutions are analogous to

Table. E.s.r. signals in M-12-crown-4 solutions

Phase (T/K)	M	g^a	$\Delta H_{pp}/G$	T_1/s	T_2/s
Frozen solution e_s^- (120)	Na ^b	2.0027(5)	0.3 ± 0.1	$(5 \pm 2) \times 10^{-6}$	$(2.2 \pm 0.7) \times 10^{-7}$
	K	2.0020(3)	2.5 ± 0.5	$(5 \pm 2) \times 10^{-5}$	$(2.6 \pm 0.5) \times 10^{-8}$
	Rb	2.0016(5)	4 ± 1	$(1.0 \pm 0.5) \times 10^{-4}$	$(1.7 \pm 0.5) \times 10^{-8}$
	Cs	2.0020(5)	3.4 ± 0.5	$(1.5 \pm 0.5) \times 10^{-5}$	$(1.5 \pm 0.5) \times 10^{-8}$
Frozen solution $M_s^+e_s^-$ (120)	Na ^b	2.0050(10)	—	—	—
	K	2.0020(3)	12 ± 1	—	—
	Rb	2.0028(10)	17 ± 1	—	—
	Cs	2.0028(10)	11 ± 1	—	—
Fluid solution (297)	Na	2.0023(3)	0.075 ± 0.01	$(2.5 \pm 1.0) \times 10^{-6}$	$(8.7 \pm 1.0) \times 10^{-7}$
	K	2.0020(3)	0.06 ± 0.01	$(5 \pm 1) \times 10^{-6}$	$(1.1 \pm 0.2) \times 10^{-6}$
	Rb	2.0023(3)	0.53 ± 0.05	$(4.5 \pm 0.5) \times 10^{-7}$	$(1.3 \pm 0.1) \times 10^{-7}$
	Cs	2.0019(3)	0.25 ± 0.05	$(3.7 \pm 1.2) \times 10^{-7}$	$(2.4 \pm 0.4) \times 10^{-7}$

^a Values in parentheses denote error in last digit. ^b Measured at 4 K.

those for K-12-crown-4. In the frozen solutions, resonances due to the electron and the ion pair $Rb_s^+e_s^-$ are observed. The linewidths of both species are independent of temperature over the range 4–200 K. The electron resonance has a Lorentzian lineshape while the shape of the ion-pair signal is between Lorentzian and Gaussian. Relaxation times for the electron resonance at 120 K are given in the Table; T_1 decreases with heating. The ion-pair signal is inhomogeneously broadened and saturates at a higher power than the solvated electron. Both signals diminish in intensity with heating.

On melting, the e.s.r. spectrum of Figure 2 collapsed to a single, homogeneous resonance as shown by the data in the Table.

(iii) *Caesium.* A frozen caesium sample displayed e.s.r. signals attributed to the electron and to an ion pair (Figure 2 and Table). A broad signal at higher g value, not shown in Figure 2, is assigned to a second ion-pair species. Relaxation times for the solvated electron were independent of temperature (Table).

The thawed sample gave e.s.r. spectra at room temperature consisting of a single Lorentzian line (Table).

(iv) *Sodium.* The e.s.r. spectra from frozen Na-12-crown-4 solutions, as shown in Figure 2, are rather different from those of the other metal-crown ether samples at low temperature. At 4 K, two distinct resonances are observed; a narrow singlet at $g = 2.0027$ and a quartet of lines centred at $g = 2.0050$ typical of hyperfine coupling to the ^{23}Na nucleus ($I = \frac{3}{2}$) in the ion pair $Na_s^+e_s^-$. The hyperfine coupling constant, $A = 3.75$ G, indicates 1.2% atomic character relative to gaseous atomic sodium ($A_{iso} = 316$ G).¹⁹ Both paramagnetic species are present in low concentration, producing rather weak e.s.r. signals. The narrow linewidth and saturation behaviour of the singlet suggest that this resonance is due to the solvated electron (Table). In the liquid phase at 297 K the e.s.r. spectrum consists of a weak Lorentzian peak (Table).

Electron Relaxation in Metal-12-Crown-4 Solutions.—(i) *Frozen solutions.* The solvated electron and ion-pair signals both have g values close to free spin as expected. The ion-pair signal is inhomogeneously broadened; this arises from hyperfine coupling to the metal nucleus at the ion-pair site.¹⁹ For the electron signal, $T_1 \neq T_2$ (Table), implying that $\omega\tau > 1$ where ω is the microwave frequency and τ the correlation time for magnetic interactions.²⁰ In contrast, we find that the electron signal in frozen potassium-hmpa solutions has $T_1 = T_2$ at 50 K.¹⁹ The correlation time may be interpreted as the inverse of the 'hopping' frequency with which the electron moves between trapping sites in the matrix; the above results indicate a lower hopping frequency in 12-crown-4. Our pulse radiolysis results in the liquid phase suggest that the electron is better solvated by

12-crown-4 than by hmpa;²¹ this is consistent with a greater trapping energy in the frozen glass and consequently a slower hopping frequency.

(ii) *Fluid solutions.* In the fluid solutions the correlation time is determined by the electron hopping rate $1/\tau_e$ and the correlation time due to diffusion in the liquid, τ_D ; equation (8). The

$$1/\tau = 1/\tau_e + 1/\tau_D \quad (8)$$

hopping rate now refers to motion of the electron among both the 'free electron' and 'ion-pair' sites. In metal-ammonia solutions both hopping and diffusion processes are significant.¹⁷ The correlation time in a liquid of viscosity η should vary with temperature as η/T .¹⁷ Viscosity data are not available for 12-crown-4 but the metal solutions were undoubtedly highly viscous at the temperature of this study which was just above their melting point. The long correlation time which this implies is consistent with the measured relaxation times T_1 and T_2 which differ by a factor of about five (Table). In metal-ammonia solutions the two relaxation times are almost equal¹⁷ indicating a considerably shorter correlation time.

Conclusions

This paper describes the use of a combination of spectroscopic techniques to demonstrate the presence of solvated electrons, metal anions, and electron-cation pairs in solutions of the alkali metals in 12-crown-4. The anions Na^- and Rb^- are spherically symmetrical with two electrons in the outer s orbital. Solvated electrons and electron-cation pairs in frozen solutions exchange slowly enough to give distinct e.s.r. signals. On melting, a single time-averaged signal is seen; electron relaxation appears to be inhibited by the high viscosity of the solutions at room temperature.

Experimental

Sample Preparation.—Solutions of alkali metals in anhydrous 12-crown-4 were prepared in Pyrex apparatus using stringent high-vacuum techniques ($\leq 10^{-4}$ mmHg). The metals were melted under vacuum and forced into capillary tubes using dried argon gas. Final metal purification was achieved by multiple distillation in the sample preparation cell. 12-Crown-4 (Aldrich) was refluxed over sodium-potassium alloy for 1 h and distilled into break-seal ampoules, then distilled again from sodium-potassium alloy in the sample preparation cell. Finally, purified solvent and metal were mixed and the resulting blue solution transferred to an optical cell, n.m.r., or e.s.r. tube as required.

Optical Absorption Spectra.—These were recorded using a Cary 17 spectrophotometer with quartz optical cells. Due to the intense absorption of the solution, a film of liquid covering the inside walls of the cell was sufficient to give a strong spectrum.

Pulse Radiolysis.—The pulse radiolysis spectrum of pure 12-crown-4 was recorded using 0.2- μ s pulses of 3-MeV electrons from a Van de Graaff generator at the Cookridge Radiation Research Centre, University of Leeds. The crown ether, purified as described above, was further distilled into a 1-cm path-length quartz radiolysis cell and degassed by three freeze-pump-thaw operations before sealing off under vacuum.

N.M.R. Spectra.—The n.m.r. spectra of ^{23}Na , $^{85,87}\text{Rb}$, and ^{133}Cs were measured using a JEOL FX-90Q pulsed Fourier-transform spectrometer; the ^{39}K nucleus was studied using a Bruker WH 400 spectrometer. Chemical shifts were measured relative to a solution of the appropriate metal chloride in D_2O .

E.S.R. Spectra.—The spectra were recorded on a Varian E-109E X-band spectrometer; g values were measured by reference to the standard 1,1-diphenyl-2-picrylhydrazyl (dpph) radical at $g = 2.0036$. Low temperatures were maintained with either an Oxford Instruments ESR-900 helium cryostat or a home-built apparatus using nitrogen gas as refrigerant.

Acknowledgements

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